# Cosmological Tests Using X-ray Observations of Clusters of Galaxies

Adam Mantz (Stanford/SLAC)

with Steve Allen, Harald Ebeling (Hawaii), David Rapetti, Robert Schmidt (Heidelberg), Glenn Morris, Andy Fabian (Cambridge), Doug Applegate, Maruša Bradač (UCSB), Evan Million, Anja von der Linden, Patrick Kelly and Alex Drlica-Wagner

Fermilab Particle Astrophysics Seminar

January 5, 2009

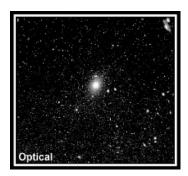


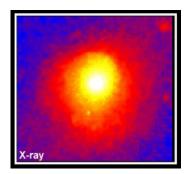




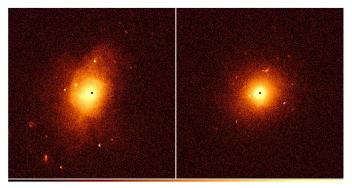
#### Why study clusters in X-rays?

- Most of the luminous matter in clusters is gas in the intracluster medium, not galaxies. In massive clusters, this gas is hot enough to radiate brightly in X-rays.
- Since X-ray luminosity depends strongly on gas density, X-ray surveys are a great way to find big clusters for cosmology.
- Primary observables (density, temperature) are closely related to the gravitational potential (total mass).



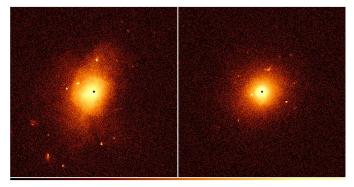


The downside: deriving interesting cluster properties from X-ray observations requires the assumption that the gas is in hydrostatic equilibrium. Simulations and mock-image analysis indicate that this biases results.

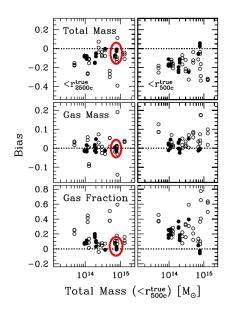


Nagai et al. 2007

The upside: the same simulations indicate that these systematics are both quantifiable and manageable.



Nagai et al. 2007

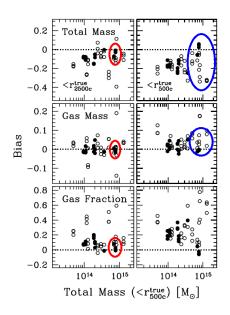


For largest relaxed clusters, we can measure at  $r_{
m 2500}$ 

- ullet  $M_{
  m gas}$  to  $\sim 1\%$  accuracy
- ullet  $M_{
  m tot}$  to few % accuracy

Bias and scatter are primarily due to non-thermal pressure support (bulk motions).

← Nagai et al. 2007 filled circles = relaxed clusters



For largest relaxed clusters, we can measure at  $r_{
m 2500}$ 

- $M_{\rm gas}$  to  $\sim 1\%$  accuracy
- $M_{
  m tot}$  to few % accuracy

Bias and scatter are primarily due to non-thermal pressure support (bulk motions).

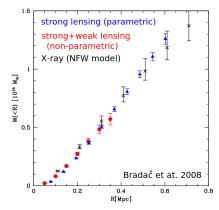
For the general population at  $r_{\rm 500}$ 

- $M_{
  m gas}$  is still recovered to few %
- $M_{
  m tot}$  is underestimated by 20–30%

← Nagai *et al.* 2007 filled circles = relaxed clusters

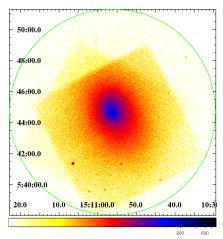
Gravitational lensing provides another handle.



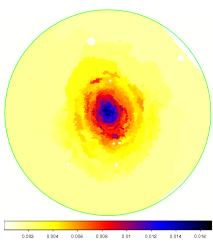


RXJ1347-1145 (z = 0.45)

#### What does relaxed mean?

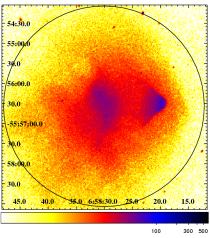


Million & Allen 2008

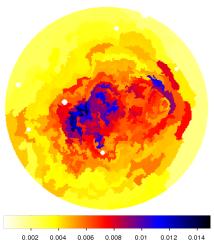


Abell 2029

#### What does unrelaxed mean?



Million & Allen 2008



1E0657-56

#### Outline

### Cluster gas-mass fraction

Measuring  $\Omega_{\rm m}$  with local observations Constraining dark energy:  $f_{\rm gas}$  as a standard ruler

#### Growth of structure

Ingredients: the mass function and cluster scaling relations Constraints on dark matter and dark energy Tests of General Relativity

#### Allen et al. 2008, MNRAS, 383, 879

(See also e.g. White & Frenk '91; Fabian '91; Briel et al. '92; White et al '93; David et al. '95; White & Fabian '95; Evrard '97; Mohr et al '99; Ettori & Fabian '99; Roussel et al. '00; Grego et al '00; Ettori et al. '03; Sanderson et al. '03; Lin et al. '03; LaRoque et al. '06; Allen et al. '02, '04.)

Fair sample hypothesis: galaxy clusters are so large that their matter content is approximately a fair sample of the matter content of the Universe (White & Frenk 1991).

For <u>relaxed</u> clusters, gas mass and total mass can be measured accurately with X-rays.

Fair sample hypothesis: galaxy clusters are so large that their matter content is approximately a fair sample of the matter content of the Universe (White & Frenk 1991).

For  $\underline{\text{relaxed}}$  clusters, gas mass and total mass can be measured accurately with X-rays.

Definitions:

$$f_{
m gas} = rac{M_{
m gas}}{M_{
m tot}}$$
  $s = rac{M_{
m star}}{M_{
m gas}} = rac{f_{
m star}}{f_{
m gas}}$   $f_{
m baryon} = f_{
m star} + f_{
m gas} = f_{
m gas}(1+s)$ 

Fair sample hypothesis: galaxy clusters are so large that their matter content is approximately a fair sample of the matter content of the Universe (White & Frenk 1991).

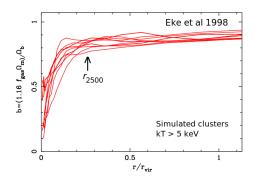
For <u>relaxed</u> clusters, gas mass and total mass can be measured accurately with X-rays.

Definitions:

$$f_{
m gas} = rac{M_{
m gas}}{M_{
m tot}}$$
  $s = rac{M_{
m star}}{M_{
m gas}} = rac{f_{
m star}}{f_{
m gas}}$   $f_{
m baryon} = f_{
m star} + f_{
m gas} = f_{
m gas}(1+s)$ 

Fair sample:  $f_{\rm baryon} = b \ \Omega_{\rm b}/\Omega_{\rm m}$ .

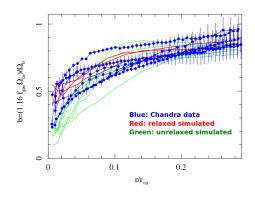
$$f_{\rm gas} = \frac{f_{
m baryon}}{1+s} = \frac{b}{1+s} \left( \frac{\Omega_{
m b}}{\Omega_{
m m}} \right)$$



#### Non-radiative simulations indicate

$$b = f_{\text{baryon}} \frac{\Omega_{\text{m}}}{\Omega_{\text{b}}} = 0.83 \pm 0.09$$

(+ 10% systematic uncertainty) at  $r_{2500}$ .

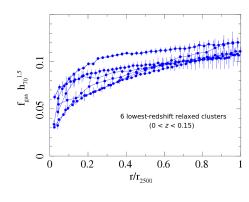


Non-radiative simulations indicate

$$b = f_{\text{baryon}} \frac{\Omega_{\text{m}}}{\Omega_{\text{b}}} = 0.83 \pm 0.09$$

(+ 10% systematic uncertainty) at  $r_{2500}$ .

Observations and simulations agree well at  $r_{2500}$ .



Non-radiative simulations indicate

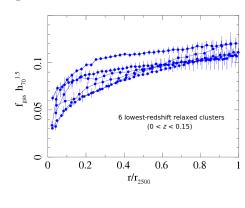
$$b = f_{\text{baryon}} \frac{\Omega_{\text{m}}}{\Omega_{\text{b}}} = 0.83 \pm 0.09$$

(+ 10% systematic uncertainty) at  $r_{2500}$ .

Observations and simulations agree well at  $r_{2500}$ .

Constant fit to the data:

$$f_{\rm gas}(r_{2500}) = (0.113 \pm 0.003) h_{70}^{-1.5}$$



Non-radiative simulations indicate

$$b = f_{\text{baryon}} \frac{\Omega_{\text{m}}}{\Omega_{\text{b}}} = 0.83 \pm 0.09$$

(+ 10% systematic uncertainty) at  $r_{2500}$ .

Observations and simulations agree well at  $r_{
m 2500}$ .

Constant fit to the data:

$$f_{\rm gas}(r_{2500}) = (0.113 \pm 0.003) h_{70}^{-1.5}$$

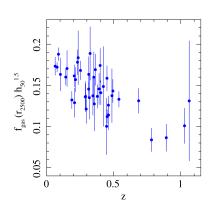
Using  $\Omega_{\rm b}h^2=0.0214\pm0.002$  (Kirkman *et al.* '03),  $h=0.72\pm0.08$  (Freedman *et al.* '01),  $s=(0.16\pm0.048)h_{70}^{1/2}$  (e.g. Lin & Mohr '04),  $b=0.83\pm0.09$  (Eke *et al.* '98 +10% systematic allowance),

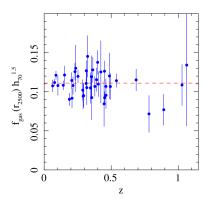
$$\Omega_{\rm m} = \frac{b \ \Omega_{\rm b}}{f_{\rm gas}(1+s)} = 0.27 \pm 0.04$$

Measured  $f_{\rm gas}$  values depend on the assumed distances to clusters as  $f_{\rm gas} \sim d^{3/2}$ . This makes the apparent  $f_{\rm gas}(z)$  dependent on cosmological parameters.

SCDM (
$$\Omega_{\rm m}=1.0,~\Omega_{\Lambda}=0.0$$
)

$$\Lambda\text{CDM}$$
 ( $\Omega_{\rm m}=0.3,~\Omega_{\Lambda}=0.7$ )

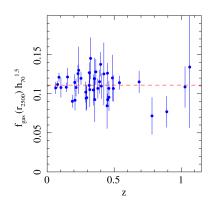




Measured  $f_{\rm gas}$  values depend on the assumed distances to clusters as  $f_{\rm gas} \sim d^{3/2}$ . This makes the apparent  $f_{\rm gas}(z)$  dependent on cosmological parameters.

Expectation: approximately constant with redshift

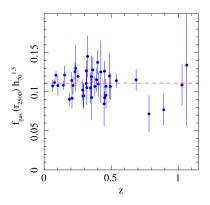
 $\Lambda\text{CDM}~(\Omega_{\rm m}=0.3,~\Omega_{\Lambda}=0.7)$ 



Measured  $f_{\rm gas}$  values depend on the assumed distances to clusters as  $f_{\rm gas} \sim d^{3/2}$ . This makes the apparent  $f_{\rm gas}(z)$  dependent on cosmological parameters.

Expectation: approximately constant with redshift

 $\Lambda\text{CDM}~(\Omega_{\rm m}=0.3,~\Omega_{\Lambda}=0.7)$ 



The full model:

$$f_{\rm gas}^{\rm ref}(z) = K \gamma \left[ \frac{b(z)}{1+s(z)} \right] \left( \frac{\Omega_{\rm b}}{\Omega_{\rm m}} \right) \left( \frac{\theta_{\rm 2500}^{\rm ref}}{\theta_{\rm 2500}^{\rm trial}} \right)^{\eta} \left[ \frac{d^{\rm ref}(z)}{d^{\rm trial}(z)} \right]^{3/2}$$

The full model:

$$f_{\rm gas}^{\rm ref}(z) = K \gamma \left[ \frac{b(z)}{1+s(z)} \right] \left( \frac{\Omega_{\rm b}}{\Omega_{\rm m}} \right) \left( \frac{\theta_{\rm 2500}^{\rm ref}}{\theta_{\rm 2500}^{\rm trial}} \right)^{\eta} \left[ \frac{d^{\rm ref}(z)}{d^{\rm trial}(z)} \right]^{3/2}$$

Conservative systematic allowances used:

- K: Instrument calibration, X-ray modeling 10% Gaussian uncertainty
- $\gamma$  : Non-thermal pressure support in gas (primarily bulk motions)  $\gamma = M_{\rm est}/M_{\rm true}$  10% uniform prior,  $1<\gamma<1.1$

The full model:

$$f_{\rm gas}^{\rm ref}(z) = K \gamma \left[ \frac{b(z)}{1+s(z)} \right] \left( \frac{\Omega_{\rm b}}{\Omega_{\rm m}} \right) \left( \frac{\theta_{\rm 2500}^{\rm ref}}{\theta_{\rm 2500}^{\rm trial}} \right)^{\eta} \left[ \frac{d^{\rm ref}(z)}{d^{\rm trial}(z)} \right]^{3/2}$$

Conservative systematic allowances used:

- b(z): Depletion factor (simulation physics, gas clumping)  $b(z) = b_0(1+\alpha_b z)$   $\pm 20\%$  uniform prior on  $b_0$   $\pm 10\%$  uniform prior on  $\alpha_b$
- s(z) : Baryonic mass in stars (observational uncertainty)  $s(z) = s_0(1+\alpha_s z) = f_{\rm star}/f_{\rm gas} \\ 30\% \ {\rm Gaussian} \ {\rm uncertainty} \ {\rm on} \ s_0 \\ \pm 20\% \ {\rm uniform} \ {\rm prior} \ {\rm on} \ \alpha_s$

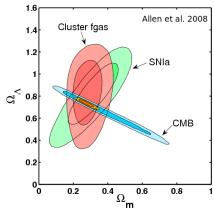
The full model:

$$f_{\rm gas}^{\rm ref}(z) = K\gamma \left[\frac{b(z)}{1+s(z)}\right] \left(\frac{\Omega_{\rm b}}{\Omega_{\rm m}}\right) \left(\frac{\theta_{\rm 2500}^{\rm ref}}{\theta_{\rm 2500}^{\rm trial}}\right)^{\eta} \left[\frac{d^{\rm ref}(z)}{d^{\rm trial}(z)}\right]^{3/2}$$

 $\Omega_{\rm b}$ : Baryon density (independent data)  $\Omega_{\rm b}h^2=0.0214\pm0.002$  (Kirkman *et al.* '03)  $h=0.72\pm0.08$  (Freedman *et al.* '01)

 $\eta$  : slope of  $f_{\rm gas}(r)$  at  $r_{2500}$  (measured) 10% Gaussian uncertainty

### $f_{\mathrm{gas}}$ : results for $\Lambda \mathsf{CDM}$ models



 $f_{\rm gas}$  alone:

$$\Omega_{\rm m} = 0.27 \pm 0.06$$
 $\Omega_{\Lambda} = 0.86 \pm 0.119$ 

Goodness of fit:  $\chi^2_{\nu}=41.5/40$ 

 $f_{\rm gas}$ : 42 clusters with standard priors

CMB: WMAP3+CBI+ACBAR + prior 
$$0.2 < h < 2$$

Supernovae: 192 from Davis '07 (ESSENCE+SNLS+HST+nearby)

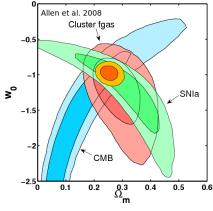
Combination does not require  $\Omega_{\rm b}h^2$ , h priors.

Combination:

$$\Omega_{\rm m} = 0.275 \pm 0.033$$

$$\Omega_{\Lambda} = 0.735 \pm 0.023$$

### $f_{\rm gas}$ : results for flat, constant-w models



 $f_{\rm gas}$  alone:

$$\Omega_{\rm m} = 0.28 \pm 0.06$$

$$w = -1.14^{+0.27}_{-0.35}$$

 $f_{
m gas}$ : 42 clusters with standard priors

$$\begin{split} \text{CMB: WMAP3+CBI+ACBAR} \\ + \text{ prior } 0.2 < h < 2 \end{split}$$

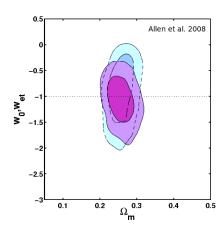
Supernovae: 192 from Davis '07 (ESSENCE+SNLS+HST+nearby)

Combination does not require  $\Omega_{\rm b}h^2$ , h priors.

#### Combination:

$$\Omega_{\rm m} = 0.253 \pm 0.021$$
 $w = -0.98 \pm 0.07$ 

### $f_{ m gas}$ : results for flat, evolving-w models



$$w(z) = \frac{w_0 z_t + w_{et} z}{z + z_t}$$

Marginalized over the transition redshift  $0.5 < 1/(1+z_t) < 0.95$ 

#### Using

 $f_{
m gas}$ : 42 clusters with standard priors

CMB: WMAP3+CBI+ACBAR + prior 
$$0.2 < h < 2$$

Supernovae: 192 from Davis '07 (ESSENCE+SNLS+HST+nearby)

Combination does not require  $\Omega_{\rm b}h^2$ , h priors.

#### Results are consistent with $\Lambda \text{CDM}$

$$w_0 = -1.05^{+0.31}_{-0.26}$$

$$w_{et} = -0.83^{+0.48}_{-0.43}$$

#### Outline

### Cluster gas-mass fraction

Measuring  $\Omega_{\rm m}$  with local observations

Constraining dark energy:  $f_{
m gas}$  as a standard ruler

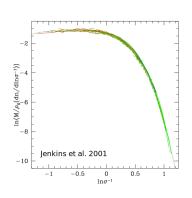
#### Growth of structure

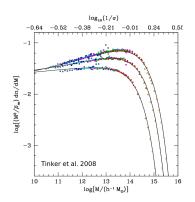
Ingredients: the mass function and cluster scaling relations Constraints on dark matter and dark energy Tests of General Relativity

Mantz *et al.* 2008, MNRAS, 387, 1179 Rapetti *et al.* 2008, arXiv:0812.2259

(See also e.g. Henry '00; Borgani et al '01; Reiprich & Böhringer '02; Seljak '02; Viana et al '02; Allen *et al.* '03; Pierpaoli *et al.* '03; Vikhlinin *et al.* '03; Schuecker et al '03; Voevodkin & Vikhlinin '04; Henry '04; Dahle '06, Henry *et al.* '08, Vikhlinin *et al.* '08)

### Growth of structure: theoretical prediction



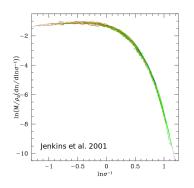


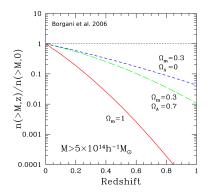
Suites of cosmological simulations predict cluster density  $\left\langle \frac{dn}{dMdz} \right\rangle$  .

The mass function can be written in universal form (within 10–20% across a range of tested cosmological models) in terms of  $\sigma(M,z)$ , where

$$\sigma^2(M,z) = \frac{1}{2\pi^2} \int_0^\infty k^2 P(k,z) |W_M(k)|^2 dk$$

### Growth of structure: theoretical prediction

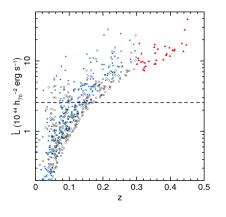




Once the normalization ( $\sigma_8$ ) is measured, the evolution of the mass function can be used to learn about dark energy.

### Growth of structure: X-ray luminosity function

Main observable: a wide-area, clean, complete cluster sample with a well understood selection function.

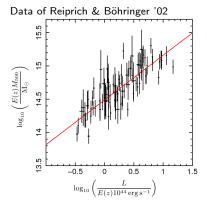


Samples based on the ROSAT All-Sky Survey:

- ▶ BCS (Ebeling et al. '98, '00) z < 0.3  $\sim 33\%$  sky coverage  $F > 4.4 \times 10^{-12} \, \mathrm{erg \, s^{-1} \, cm^{-2}}$
- ▶ REFLEX (Böhringer et al. '04) z < 0.3  $\sim 33\% \text{ sky coverage}$   $F > 3.0 \times 10^{-12} \, \mathrm{erg \, s^{-1} \, cm^{-2}}$
- ► MACS (Ebeling et al. '01, '07) 0.3 < z < 0.7  $\sim 55\%$  sky coverage  $F > 2.0 \times 10^{-12} \, \mathrm{erg \, s^{-1} \, cm^{-2}}$

Luminosity cut at  $2.55 \times 10^{44} h_{70}^{-2} \, \mathrm{erg \, s^{-1}}$  leaves  $78 + 130 + {34 \over 34} = 242$  massive clusters.

### Growth of structure: scaling relation



Self-similarity suggests a power-law model

$$Y = \alpha + \beta X_1$$

where

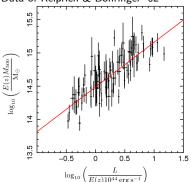
$$Y = \log_{10} \left( \frac{E(z)M_{500}}{M_{\odot}} \right)$$

$$X_{1} = \log_{10} \left( \frac{L}{E(z)10^{44} \,\mathrm{erg}\,\mathrm{s}^{-1}} \right)$$

and the factors of  $E(z)=H(z)/H_0$  are due to the evolution in  $r_{500}$ .

### Growth of structure: scaling relation

Data of Reiprich & Böhringer '02



Self-similarity suggests a power-law model

$$Y = \alpha + \beta X_1$$

where

$$Y = \log_{10} \left( \frac{E(z)M_{500}}{M_{\odot}} \right)$$

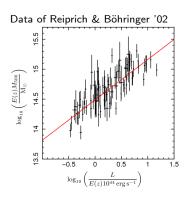
$$X_{1} = \log_{10} \left( \frac{L}{E(z)10^{44} \,\mathrm{erg \, s^{-1}}} \right)$$

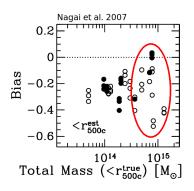
and the factors of  $E(z) = H(z)/H_0$  are due to the evolution in  $r_{500}$ .

Marginalize over possible deviations from self-similarity:

$$Y = \alpha + \beta X_1 + \gamma X_2$$
  
$$X_2 = \log_{10}(1+z)$$

### Growth of structure: scaling relation



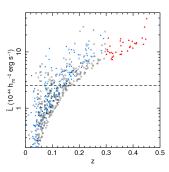


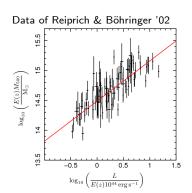
M, L from X-ray observations (Reiprich & Böhringer '01). Departures from hydrostatic equilibrium introduce a bias in the estimates of  $r_{500}$ ,  $M_{500}$ .

Based on simulations, marginalize over bias  $-25(\pm 5)\%$  and scatter  $\pm 15(\pm 3)\%$ .

A problem, but major improvements are possible...in preparation.

#### **BCS**+REFLEX+MACS



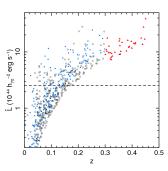


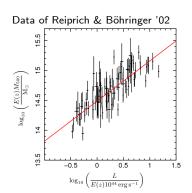
Compare the survey data (redshifts, fluxes) with the expectation:

$$\left\langle \frac{dN(z,\hat{L})}{dzd\hat{L}} \right\rangle = P_{\rm sel}(z,\hat{L}) \int_0^\infty dL \ P(\hat{L}|L) \int_0^\infty dM \ P(L|M) \left\langle \frac{dn(z,M)}{dM} \right\rangle \frac{dV}{dz}$$

selection function

#### **BCS**+REFLEX+MACS



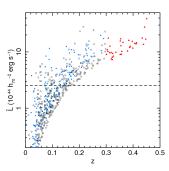


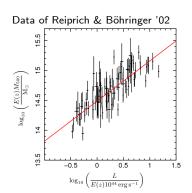
Compare the survey data (redshifts, fluxes) with the expectation:

$$\left\langle \frac{dN(z,\hat{L})}{dzd\hat{L}} \right\rangle = P_{\rm sel}(z,\hat{L}) \int_0^\infty dL \ P(\hat{L}|L) \int_0^\infty dM \ P(L|M) \left\langle \frac{dn(z,M)}{dM} \right\rangle \frac{dV}{dz}$$

flux measurement error

#### **BCS**+REFLEX+MACS



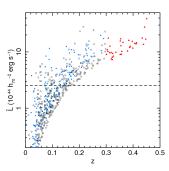


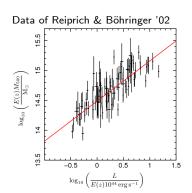
Compare the survey data (redshifts, fluxes) with the expectation:

$$\left\langle \frac{dN(z,\hat{L})}{dzd\hat{L}} \right\rangle = P_{\rm sel}(z,\hat{L}) \int_0^\infty dL \ P(\hat{L}|L) \int_0^\infty dM \ P(L|M) \left\langle \frac{dn(z,M)}{dM} \right\rangle \frac{dV}{dz}$$

mass-luminosity intrinsic scatter

#### **BCS**+REFLEX+MACS





Compare the survey data (redshifts, fluxes) with the expectation:

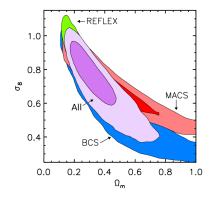
$$\left\langle \frac{dN(z,\hat{L})}{dzd\hat{L}} \right\rangle = P_{\rm sel}(z,\hat{L}) \int_0^\infty dL \ P(\hat{L}|L) \int_0^\infty dM \ P(L|M) \left\langle \frac{dn(z,M)}{dM} \right\rangle \frac{dV}{dz}$$

mass function and volume element

### Growth of structure: priors and systematic allowances

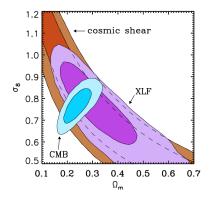
Cosmological parameters		
Hubble constant, $h$	$0.72 \pm 0.08$	Hubble Key
Baryon density, $\Omega_{ m b} h^2$	$0.0214 \pm 0.002$	BBN
Mass function		
normalization	$\pm 20\%$ Gaussian	
Mass-luminosity relation		
non-similar evolution	$\pm 20\%$ uniform	
scatter evolution	$\pm 30\%$ uniform	
mass bias and scatter	$\pm 20\%$ Gaussian	

#### Growth of structure: results for flat $\Lambda$ CDM models



Results from the 3 cluster samples individually are consistent with one another, and with most previous work using clusters.

#### Growth of structure: results for flat $\Lambda$ CDM models



Results agree with independent cosmological data.

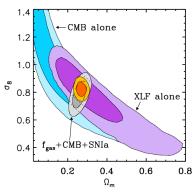
XLF: 242 clusters, z < 0.7

$$\Omega_{\rm m} = 0.28^{+0.11}_{-0.07} 
\sigma_8 = 0.78^{+0.11}_{-0.13}$$

CMB: WMAP3 data

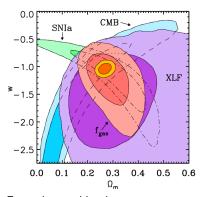
Cosmic shear: CFHTLS Wide, linear regime (Fu et al. '08)

#### Growth of structure: results for flat constant-w models



#### From the XLF:

$$\Omega_{\rm m} = 0.24^{+0.15}_{-0.07} 
\sigma_8 = 0.85^{+0.13}_{-0.26} 
w = -1.4^{+0.4}_{-0.7}$$

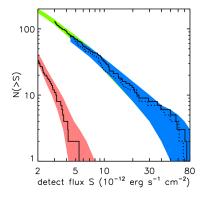


#### From the combination:

$$\Omega_{\rm m} = 0.269 \pm 0.016 \quad (0.258 \pm 0.022)$$
 $\sigma_8 = 0.82 \pm 0.03 \quad (0.79 \pm 0.06)$ 
 $w = -1.02 \pm 0.06 \quad (-0.99 \pm 0.07)$ 

Combined results are consistent with  $\Lambda$ CDM.

### Growth of structure: goodness of fit



Observed number as a function of flux limit compared with predictions the  $\Lambda \text{CDM}$  best fit for each survey.

### Testing general relativity with the growth of structure

Growth of density perturbations

$$\delta = (\rho_{\rm m} - \bar{\rho}_{\rm m})/\bar{\rho}_{\rm m}$$
 in GR:

$$\ddot{\delta} + 2\frac{\dot{a}}{a}\dot{\delta} = 4G\pi\rho_{\rm m}\delta$$

Instead, parametrize through

$$\frac{d\delta}{da} = \frac{\delta}{a} \Omega_{\rm m}(a)^{\gamma}$$

with  $\gamma \sim 0.55$  accurately reproducing GR.

The growth function in this model

- matches GR at early times
- has the same scale dependence as GR
- is allowed to have a different time dependence

On smaller scales, gravity is unmodified.

### Testing general relativity with the growth of structure

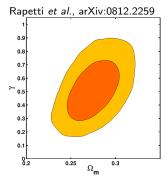
Growth of density perturbations  $\delta = (\rho_m - \bar{\rho}_m)/\bar{\rho}_m$  in GR:

$$\ddot{\delta} + 2\frac{\dot{a}}{a}\dot{\delta} = 4G\pi\rho_{\rm m}\delta$$

Instead, parametrize through

$$\frac{d\delta}{da} = \frac{\delta}{a} \Omega_{\rm m}(a)^{\gamma}$$

with  $\gamma \sim 0.55$  accurately reproducing GR.



 $\mathsf{XLF} + \mathsf{WMAP5} + f_{\mathsf{gas}} + \mathsf{snIa}(\mathsf{Union})$ 

$$\Lambda {
m CDM}$$
  $\gamma = 0.51^{+0.16}_{-0.15}$   $w {
m CDM}^*$   $\gamma = 0.44^{+0.17}_{-0.15}$  non-flat  $\Lambda {
m CDM}$   $\gamma = 0.51^{+0.19}_{-0.14}$ 

 $<sup>^*</sup>w$  is used only to parametrize the expansion history in this model

#### Conclusions

•  $f_{\rm gas}(z)$  data for largest relaxed clusters  $\to$  tight constraints on  $\Omega_{\rm m}$ ,  $\Omega_{\Lambda}$  and w through absolute distance measurements.

$$\Omega_{\rm m} = 0.27 \pm 0.06$$
  $\Omega_{\Lambda} = 0.86 \pm 0.119$   $\left(w = -1.14^{+0.27}_{-0.35}\right)$ 

• Growth of X-ray luminous clusters spanning  $0 < z < 0.7 \to$  independent constraints on  $\Omega_{\rm m}$ ,  $\sigma_8$  and w.

$$\Omega_{\rm m} = 0.28^{+0.11}_{-0.07} \qquad \sigma_8 = 0.78^{+0.11}_{-0.13} \qquad \left(w = -1.4^{+0.4}_{-0.7}\right)$$

ightharpoonup Combination of  $f_{gas}$ , XLF, CMB and snla data

$$\Omega_{\rm m} = 0.269 \pm 0.016$$
  $\sigma_8 = 0.82 \pm 0.03$   $w = -1.02 \pm 0.06$ 

Combination applied to tests of General Relativity

$$\Omega_{\rm m} = 0.27 \pm 0.02 \qquad \sigma_8 = 0.82 \pm 0.05 \qquad \gamma = 0.51 \pm 0.15$$

This year: new X-ray and lensing data should provide improvement in both  $f_{
m gas}$  and XLF results

Next few years: new  $f_{\rm gas}$  targets and cluster samples from SZ, X-ray, optical surveys

